From Table 3 it can be noted that each atom in the structure has a coordination atmosphere of 14 near neighbors with the coordinated atoms forming nearly regular rhombic dodecahedra with one of the major axes pointing in the [$\overline{160}$] direction. The volume contraction associated with phase formation from the solid elemental components has been computed as 3%. The structure shows pronounced layering of the atoms in the (001), (610), ($\overline{830}$), (7 $\overline{12}$), ($\overline{121}$), and (614) planes. This layering may well be associated with low activation energies for dislocation glide since it was observed qualitatively that Li₉Al₄, in contrast to Na₂Tl, could suffer appreciable plastic strain without fracture.

There is a formal relationship between the structure of Li_9Al_4 and the tetragonal MoSi_2 -type structure in the sense that, if one considers only atomic loci and not species, atomic arrays in the two structures are closely comparable. Both structures are formed by alternating atomic layers normal to the *c* axis in the case of Li_9Al_4 and normal to the *a* or *b* axes in the case of MoSi_2 . The relationship can be observed by comparing Figs. 3 and 4. This relationship between the two structures may simply reflect the fact that geometric constraints govern the efficient filling of space, and both Li_9Al_4 and $MoSi_2$ are efficiently packed structures.

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The Crystal Structure of Tetracyanocyclobutane

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cis,trans,cis-1,2,3,4-Tetracyanocyclobutane, C₈H₄N₄, crystallizes in the monoclinic system; the space group is $P2_1/n$, with a=8.610, b=6.193, c=6.983 Å (all ± 0.005 Å), $\beta=98.4\pm0.1^{\circ}$. There are two molecules per unit cell. The crystal structure has been determined from counter data with use of the symbolic addition phase determining procedure. The cyclobutane ring is planar and rectangular with C-C edges of 1.561 ± 0.003 and 1.547 ± 0.003 Å. Bond lengths in the C-C=N 'legs', corrected for the effects of thermal motion, are C-C, 1.452 ± 0.005 Å and C=N, 1.158 ± 0.008 Å.

Introduction

The configuration of the highly strained cyclobutane ring in organic compounds has long been of considerable interest, but until recently only a few serious efforts had been made to investigate it by the methods of X-ray diffraction. These are summarized in Table 1. It appears that the cyclobutane ring is puckered in some compounds and planar in others. Additional studies are clearly needed to establish the conditions favoring one configuration over the other, as well as to determine bond lengths with maximum precision.

Tetracyanocyclobutane (TCCB) was synthesized by Griffin, Basinsky & Peterson (1962) by the solid state photodimerization of fumaronitrile. Chemical and physical tests indicated that the product of the reaction was probably the *cis*, *trans*, *cis* isomer (Fig. 1).

Both TCCB and fumaronitrile crystallize in the monoclinic system, space group $P2_1/n$, with two mol-



Fig. 1. Dimerization of fumaronitrile to form tetracyanocyclobutane.

^{*} Taken from a thesis submitted to the Faculty of the Polytechnic Institute of Brooklyn in partial fulfillment of the requirements for the degree of Master of Science.

ecules per unit cell. The crystal structure of fumaronitrile was determined by Orsa (1960). It seemed likely that the analysis of the crystal structure of TCCB would not only establish the configuration of its cyclobutane ring, but might reveal structural relationships between the molecular arrangements in crystals of the monomer and dimer.

Experimental

Single crystals of TCCB suitable for X-ray diffraction analysis were grown by slow evaporation of solutions in benzene and in carbon tetrachloride. Both solvents yielded the same product. A small irregularly shaped crystal ($0.1 \times 0.2 \times 0.3$ mm) was selected for investiga-

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	Dihedral angle in ring	C-C bond length(s) in ring	Reference
Tetraphenylcyclobutane	180°	1.555 ± 0.02 Å 1.585	Dunitz (1949)
	180	1.566 ± 0.015 1.573	Margulis (1965 <i>b</i>)*
Octachlorocyclobutane	158	1·58 1·60	Owen & Hoard (1951)
Photodimer of cyclopentenone	180	1·54 1·60	Margulis (1965 <i>a</i>)
cis-1,2-Dibromo-1,2-dimethoxycarbonylcyclobutane	150	$ \begin{array}{r} 1 \cdot 52 \\ 1 \cdot 58 \\ 1 \cdot 58 \\ 1 \cdot 58 \\ 1 \cdot 55 \end{array} $	Karle, Karle & Britts (1966)
trans-1,2-Dibromo-1,2-dimethoxycarbonylcyclobutane	153	${ \begin{array}{c} 1 \cdot 52 \\ 1 \cdot 57 \\ 1 \cdot 58 \end{array} \pm 0 \cdot 027 \\ 1 \cdot 58 \end{array} }$	Karle, Karle & Britts (1966)
Anemonin	152	$ \frac{1 \cdot 530 \pm 0.01}{1 \cdot 545} \\ \frac{1 \cdot 541 \pm 0.017}{1 \cdot 532} $	Karle & Karle (1966)
Octahydroxycyclobutane	180	1.562 average	Bock (1967)
trans-1,3-Cyclobutane-dicarboxylic acid	180	1·552 1·567 <u>+</u> 0·006	Margulis & Fisher, M.S. (1967)
cis, trans, cis-1,2,3,4-Tetracyanocyclobutane	180	1·547 ± 0·003 1·561	This study

* Refinement of Dunitz's data.



Fig. 2. Thermal motions of atoms in tetracyanocyclobutane.

Crystal system Space group	$\begin{array}{c} \text{TCCB} \\ \text{Monoclinic} \\ P2_1/n \end{array}$	Fumaronitrile Monoclinic P2 ₁ /n
a	8.610	3.89 Å
<i>b</i>	$6.193 \} \pm 0.005 A$ 6.983 1	5.60 10.04
β	$98.4 \pm 0.1^{\circ}$	102·5°
Ζ	2	2
QX-ray	1.408 g.cm^{-3}	1.21 g.cm ⁻³
Qmeas.	1.384	1.196 ± 0.005

Table 2. Crystal data

tion. Crystal data for TCCB and fumaronitrile are listed in Table 2.

Intensity data were obtained with a single-crystal diffractometer using filtered Cu Ka radiation and pulse height analysis to reduce unwanted background. Reflection intensities were measured to $2\theta = 150^{\circ}$ by the stationary-crystal stationary-counter method. Six hundred and fifty-one of a total of 765 reflections accessible to Cu $K\alpha$ radiation were detected and measured. General background was measured as a function of reflection angle and subtracted from each reflection intensity. In addition, forty-seven reflections were remeasured with the use of ' θ -2 θ ' scans to obtain 'integrated' intensities. The ratios of these intensities to the corresponding 'peak' intensities were plotted as functions of reflection angles to convert 'peak' to 'integrated' intensities. The usual geometrical corrections were then applied to the latter.

Structure determination

Although the structure was clearly a simple one, we experienced considerable difficulty in interpreting 'sharpened' and 'unsharpened' Patterson maps computed from the three-dimensional data. The causes of these difficulties became apparent when the structure determination was completed and will be discussed below.

At this point we decided to use the symbolic addition procedure of Karle & Karle (1963) to determine the signs of the structure factors. Using a program described by Bednowitz (1966), E factors were computed from the experimental intensities. Six reflections with large E factors and many interactions were selected to initiate the phase determining procedure. Three of the six were assigned positive signs to fix the origin of the unit cell; the three others were given symbolic signs A, B and C (Table 3).

Table 3.	Phase	assign	iments j	for	the
symb	olic ad	dition	proced	ure	

			Number
h, k, l	E	Signs	of interactions
10,0,4	3.34	В	41
4,3,4	3.11	+	44
3,2,4	3.01	+	51
1,2,0	2.96	A	58
4,3,6	2.80	С	45
2,2,3	2.44	+	43

Symbolic signs of approximately 270 reflections with |E| > 0.90 were calculated, including 29 h0l terms. With the use of these 29 terms, eight projections of electron density onto the *ac* plane were computed with the von Eller (1955) photosommateur. These projections were based on the eight phase combinations of the three symbolic signs, A reasonable peak arrangement was detected in one of the projections and the values of the symbolic signs based on that arrangement were used to compute (*hk*0) and (0*kl*) electron density maps. These confirmed that the correct choice had been made.

Coordinates obtained from these electron density maps served as the starting positions for least-squares

Table 4. Atomic position and thermal parameters

Estimated standard deviations $\times 10^4$ in parentheses.

	A. Position	parameters				
		x		<i>y</i>	Z	
	C(1)	0.0362 (2)	0.08	353 (3)	-0.1245(3)	
	C(2)	0.1499 (3)	0.25	579 (3)	-0.1217(3)	
	C(3)	0.0921 (2)	-0.12	211 (3)	-0.0059(3)	
	C(4)	0.2504 (2)	0·11	51 (4)	0.0952(3)	
	N(1)	0.2421(3)	0.38	384 (3)	-0.1150(4)	
	N(2)	0.3733 (2)	-0.10)55 (4)	0.1753 (3)	
	H(1)	-0.0069 (20)	0.04	435 (36)	-0.2528(29)	
	H(2)	0.0758 (22)	-0.23	541 (34)	-0.0830 (28)	
B. Ther (base	rmal parameter ed on $b_{11}h^2 + b$	rs (×10 ⁴) $b_{22}k^2 + b_{33}l^2 + 2$	$b_{12}hk + 2b_{13}$	$hl + 2b_{23}kl$)		
	b_{11}	b22	b33	<i>b</i> ₁₂	b ₁₃	b23
C(1)	79 (3)	155 (6)	118 (4)	8 (3)	5 (3)	8 (3)
C(2)	99 (3)	177 (6)	170 (5)	15 (3)	40 (3)	25 (3)
C(3)	69 (3)	137 (5)	135 (4)	4 (3)	4 (3)	-5(3)
C(4)	77 (3)	230 (6)	181 (5)	12 (3)	8 (3)	53 (4)
N(1)	145 (4)	206 (6)	334 (6)	-23 (3)	105 (4)	18 (4)
N(2)	92 (3)	508 (10)	270 (5)	2 (4)	- 34 (3)	132 (6)
H(1)	B = 1.0 (0.4)	4)				
H(2)	B = 0.5 (0.4)	4)				

••	K L	FOBS	FCAL	HKL	FOBS	PCAL	H	K L	FOBS	PCAL	H	ΚĽ	FOBS	FCAL	H	K L	POBS	PCAL	H	K L	POBS	PCAL
1	00001111111112222222233333334444445555566666777 000000001111111111111111122222222222	1340977455173043381777655144087666622441934087014811115534564174866 819162736784387137999267724493147092588824889324993469752366612975236665529417150584342712	134207310508319447719959364183958044473475911292191512145147380391863138868152744529241957544431255333555518311788256244399718081445997461310483299312	~~~~~~~~~~~~~~~~~~~~~~~~~~~~~~~~~~~~~	233537293226773359991615722875253224395 101739934413703022766122499170081431133726612240917394011227049214236116551261269362781018101739447218452184154	2833437301895814213338171815232827287402203884 12078122331732433148866352253658131958814213338878812781017117446822253658	3	1777 0000001111111111111111111111111111	11062 267331830929163719268239991875386 844471318001724341 79862667144475132926365114442060395522191411733254 33020175190271466482211155668981125	1152 297312733062661595021788877373882423521711334648671835910096848734596673499642300535402682707314575392054202142180241 913251207492214524192346378301123	5	22222222222222222222222222222222222222	22180365248181886666628027745243015112266026831921624518627100245883418270148287137 2810405716901153118244197502638024483973330803443997333080755569733856653101610	3121822771193397647647832702718551425014029377781132081743286925983249933501834281411287235 291039092710331131372311897530141427422112624915188893322694944055249968	6	44444555555555555555555555555555555555	11342344042910522602619184362519 324695133537902152763692812873684188712744388309331604333420416889844522190640113424 11799744082327306666438934471151074408232730666645889841889113222190640113424 1179744082327306666458898412822190640113424 1179744082327306666458898412822190640113424 1179744082327306666458841889113222190640113424 1179744082327306666458884188911322190640113424 117991440823273066664588418891132218000000000000000000000000000000000	109543141015211114825759171312616 3345983051542019318152703102162433503663845386230422872637173122511148257591715132616 334598305154201931815267391284333366384538323042287265113422413251854504413232 2019991355288127728235244733324191252266631377	9	v2355555544444444444444445555555 0000001111111111	2645755413777146543008994151268 15673148599989864990215149368981277329390987 61913626693360747722544616254296630252395531425	2639834774207833937308222016777 7690004479869463710117159859777322335998822 601611154454572918922045171572 10360627214437166152

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refinements of the structure. The BNL* version of the Busing-Levy, full-matrix least-squares program was used. This minimizes:

$$\sum_{hkl} W_{hkl} [(F_{obs})_{hkl} - (F_{calc})_{hkl}]^2.$$

Initially 109 reflections, recorded at relatively low Bragg angles, were used to refine the positions of the carbon and nitrogen atoms. After a few cycles

$$R = \sum_{hkl} \left[\frac{|(F_{obs})_{hkl} - [scalefactor] (F_{calc})_{hkl}|}{|(F_{obs})_{hkl}|} \right]$$

decreased to 0.18. All the experimental data were then included in the least-squares calculations. Anisotropic temperature factors were computed for the carbon and nitrogen atoms. The positions and isotropic temperature factors of the hydrogen atoms, which were initially located on the basis of geometrical considerations, were also refined.

Late in the refinement process, it became clear that the measured intensities of ten reflections were grossly in error and the corresponding reflections were discarded. The final value of R for the remaining 641 reflections was 0.054.

In the final refinement cycle the changes of atomic parameters, excluding those of hydrogen atoms, were less than two per cent of their estimated standard deviations. A difference electron density map was then computed which confirmed the locations of the hydrogen atoms. The largest background peaks were less than one third the height of the hydrogen peaks on the difference maps. The final atomic positions and thermal parameters are listed in Table 4. Observed and calculated structure factors are shown in Table 5.

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Three weighting schemes were used in the final stage of the least-squares refinement process, *i.e.* $\sigma = \text{con-}$ stant; $\sigma \propto F$, and a more elaborate scheme in which σ was evaluated by taking explicit account of statistical counting errors in peak intensities, background, and errors made in relating peak heights to integrated intensities. Bond lengths and angles obtained from all three weighting schemes were within two estimated standard deviations of one another. Results reported in Table 6 are based on the third of the schemes mentioned above.

Discussion

The cyclobutane ring in TCCB is planar and rectangular with bonds of 1.561 Å between *cis* C-C \equiv N 'legs' and 1.547 Å between *trans* C-C \equiv N 'legs' (each ± 0.003 Å).

The difference between the two bond lengths equals five estimated standard deviations and is therefore statistically significant. These results agree well with previous studies involving compounds with planar cyclobutane rings (Table 1). Our efforts to establish the $C \equiv N$ bond lengths with high accuracy were relatively unsuccessful. The values obtained from the least squares analysis (1.123 and 1.129 Å) were unexpectedly low, apparently owing to the large anisotropic thermal motions of these atoms (Fig. 2).

To correct these bond lengths for the effects of thermal motion, individual atoms were assumed to 'ride' on their neighbors (Busing & Levy, 1964). This led to calculated increases of 0.007 Å in the average C-C bond (from 1.445 to 1.452 Å) and 0.032 Å in the average C \equiv N bond (from 1.125 to 1.158 Å). The latter value is close to the expected one and it is tempting to view this as a confirmation of the assumed model

C(1)–C(3) C(1)–C(3)' C(1)–C(2) C(3)–C(4) C(1)–C(2)*	1·561 Å 1·547 1·448 1·441 1·453	C(3)'-C(1)-C(3) C(1)'-C(3)-C(1) C(1) -C(2)-N(1) C(3) -C(4)-N(2) C(3) -C(1)-C(2)	89·8° 90·2 177·6 178·3 116·2
C(3)-C(4)* C(2)-N(1) C(4)-N(2) C(2)-N(1)* C(4)-N(2)*	1-451 1-129 1-123 1-155 1-162	C(3)'-C(1)-C(2) C(1) -C(3)-C(4) C(1)'-C(3)-C(4) C(3) -C(1)-H(1) C(3)'-C(1)-H(1)	$ \begin{array}{r} 115.2 \\ 115.8 \\ 114.4 \\ 108.9 \pm 1.3 \\ 112.3 \pm 1.0 \end{array} $
C(1)-H(1) C(3)-H(2)	$\begin{array}{c} 0.95 \pm 0.02 \\ 0.98 \pm 0.02 \end{array}$	C(1) -C(3)-H(2) C(1)'-C(3)-H(2) C(2) -C(1)-H(1) C(4) -C(3)-H(2)	$\begin{array}{c} 112 \cdot 8 \pm 1 \cdot 1 \\ 112 \cdot 7 \pm 1 \cdot 1 \\ 112 \cdot 4 \pm 1 \cdot 2 \\ 109 \cdot 9 \pm 1 \cdot 1 \end{array}$
	Other intramolecul	ar distances	
C(1)-C(1)' C(3)-C(3)'	2·201 Å 2·194	C(2) -C(4) N(1)-N(2)	2·828 Å 3·751

Table 6. Bond lengths and angles in tetracyanocyclobutane

* Corrected for thermal motion.

The estimated standard deviations of bonds, involving non-hydrogen atoms only, equal ± 0.003 Å; the corresponding e.s.d. of angles involving these atoms equal $\pm 0.2^{\circ}$.

of the thermal motion, but this can only be unambiguously established on the basis of low temperature work.

Both $C-C \equiv N$ angles in TCCB are approximately 178°. A similar non-linearity of the $C-C \equiv N$ 'legs' has been noted in tetracyanoquinodimethane by Hanson (1965) and *cis*-1,2,3-tricyanocyclopropane by Hartman

& Hirshfeld (1966). Hartman & Hirshfeld (1966) and Dunitz & Ibers (1967) have discussed possible causes of this effect.

The closest intermolecular approaches are listed in Table 7 and are illustrated in Fig. 3. All involve at least one nitrogen atom. N1(1), for example is less than

Table 7. Closest intermolecular approaches in TCCB (Å)

For $H \cdots X$ less than 3.0 Å and for $X \cdots X'$ less than 3.5 Å, where X is C or N. The second number in the atom symbol designates:



Fig. 3. The crystal structure of tetracyanocyclobutane as viewed along the a axis.



Fig.4. The crystal structure of tetracyanocyclobutane as viewed along the b axis.

3.5 Å from ten adjacent atoms and N2(1) is less than 3.5 Å from seven atoms. The molecular packing is also shown in Fig.4. It can be seen that the molecules are most closely packed in layers approximately parallel to the (101) and (101) planes, leaving large open channels along $\frac{1}{2}$, y, 0 and 0, y, $\frac{1}{2}$.

We have mentioned that TCCB was synthesized by the solid state photodimerization of fumaronitrile. Griffin *et al.* (1962) have pointed out that ... 'the orientation (of fumaronitrile molecules in one crystal) is completely consistent with the formation of a dimer having the structure of tetracyanocyclobutane providing bond formation takes place between nearest neighbor molecules and no major shifts in relative positions of the atoms take place during the reaction'...

Our analysis of the two crystal structures indicates, in agreement with Griffin, that the arrangement of molecules in fumaronitrile, in parallel layers 3.16 Å apart, does indeed favor the formation of molecules of the dimer (TCCB) from pairs of nearest molecular neighbors. We are unable, however, to detect any significant influence of the crystal structure of fumaronitrile on the crystal structure of TCCB.

Analysis of the completed structure revealed some of the causes of our difficulty in interpreting the Patterson maps. A large number of peaks, unrelated by symmetry, are separated by $y=\frac{1}{2}$. This factor, plus the large anisotropic thermal motions of the atoms in the $C-C \equiv N$ 'legs', made interpretation of the Patterson maps difficult. By contrast the symbolic addition method proceeded smoothly and rapidly.

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